This article was downloaded by: [University of California, San Diego]

On: 11 August 2012, At: 10:35 Publisher: Taylor & Francis

Informa Ltd Registered in England and Wales Registered Number: 1072954 Registered office: Mortimer House, 37-41 Mortimer Street, London W1T 3JH,

UK



Molecular Crystals and Liquid Crystals

Publication details, including instructions for authors and subscription information: http://www.tandfonline.com/loi/qmcl20

Self-Supporting Liquid Crystal Composite

Julie Cleaver $^{\rm a}$, Doris Vollmer $^{\rm a}$, Jason Crain $^{\rm a}$ & Wilson Poon $^{\rm a}$

^a School of Physics & Astronomy & COSMIC*, University of Edinburgh, Edinburgh

Version of record first published: 18 Oct 2010

To cite this article: Julie Cleaver, Doris Vollmer, Jason Crain & Wilson Poon (2004): Self-Supporting Liquid Crystal Composite, Molecular Crystals and Liquid Crystals, 409:1, 59-68

To link to this article: http://dx.doi.org/10.1080/15421400490435666

PLEASE SCROLL DOWN FOR ARTICLE

Full terms and conditions of use: http://www.tandfonline.com/page/terms-and-conditions

This article may be used for research, teaching, and private study purposes. Any substantial or systematic reproduction, redistribution, reselling, loan, sub-licensing, systematic supply, or distribution in any form to anyone is expressly forbidden.

The publisher does not give any warranty express or implied or make any representation that the contents will be complete or accurate or up to date. The accuracy of any instructions, formulae, and drug doses should be independently verified with primary sources. The publisher shall not be liable for any loss, actions, claims, proceedings, demand, or costs or damages

whatsoever or howsoever caused arising directly or indirectly in connection with or arising out of the use of this material.

Mol. Cryst. Liq. Cryst., Vol. 409, pp. 59–68, 2004 Copyright © Taylor & Francis Inc.

ISSN: 1542-1406 print/1563-5287 online DOI: 10.1080/15421400490435666



SELF-SUPPORTING LIQUID CRYSTAL COMPOSITE

Julie Cleaver, Doris Vollmer, Jason Crain, and Wilson Poon School of Physics & Astronomy & COSMIC*, University of Edinburgh, James Clerk Maxwell Building, Mayfield Road, Edinburgh, EH9 3JZ

> Colloid-liquid crystal composites are a novel type of soft condensed matter, formed from dispersions of colloidal particles in liquid crystalline hosts. Here, we investigate the kinetics of network formation in mixtures of the thermotropic liquid crystal 5CB and nearly hard sphere colloids (polymethylmethacrylate particles), occurring as the liquid crystal undergoes a transition from the isotropic to the nematic phase. In the isotropic phase the particles (volume fraction <0.2) are homogeneously distributed. While cooling through the isotropic-nematic phase transition the particles are expelled by the liquid crystal, forming a three dimensional network structure. Using time-resolved laser scanning confocal microscopy (LSCM) we image the dynamics of the network formation at single particle resolution, for the first time. In the isotropic phase the particles perform Brownian motion, whereas in the network they remain frozen in position. We observe a broad distribution of the size of the liquid crystal domains of up to ~50 micrometers. Heating shows although the network breaks up clusters of particles remain. DSC measurements show that the particles do not change T_{Nb} suggesting that the particles do not alter the bulk properties of the liquid crystal. These observations provide an important step towards characterising the material properties of this novel composite.

Keywords: fluorescence microscopy; liquid crystal composites; mechanical properties; microstructure; network formation

*COSMIC: Collaborative Optical Spectroscopy Micromanipulation and Imaging Center.

Julie Cleaver and Doris Vollmer are funded by the Engineering and Physical Sciences
Research Council (EPSRC) and the European Commission under the Marie Curie program
respectively. We would like to thank Dr. Andy Schofield for preparing the PMMA particles.

Address correspondence to Julie Cleaver, School of Physics & Astronomy & COSMIC University of Edinburgh, James Clerk Maxwell Building, Mayfield Rd, Edinburgh, EH9 3JZ United Kingdom.

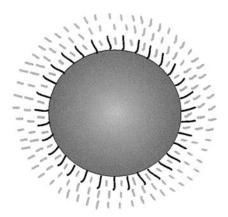


FIGURE 1 Schematic showing LC alignment around colloid particlel in the nematic phase. The LC molecules and stabilizing hairs are of lengths 2 nm and 10 nm respectively. The particle radius ranges from 100–700 nm.

1. INTRODUCTION

It has recently been discovered that a mixture of solid colloidal particles and nematic liquid crystal combines the switching properties of liquid crystals [1,2,3] and the mechanical rigidity of a wax-like soft solid [4,5,6]. Separately the two components have been studied extensively, however the area in which they are combined has entertained very little research.

A poor understanding of these mixtures exists partially because previous attempts have observed macroscopic phase separation between the particles and the liquid crystal. This is not surprising: the ordering of liquid crystal molecules imposed by the boundary condition at particle surfaces (see Figure 1) means necessarily that particles will be associated with defects in a nematic phase. These defects cost energy, and are therefore unfavourable.

The consequence is that a nematic liquid crystal will tend to exclude particles. Therefore, much of the literature on particles in liquid crystal is directed towards the study of single particles (or liquid droplets), and the structure of defects surrounding them [7,8].

There is also interest in how these defects induce anisotropic interactions and recently conditions have been found under which such mixtures remains stable. Poulin *et al.* [9] dispersed micrometer sized surfactant covered water droplets in 5CB. The colloidal droplets formed chains if the LC was in the nematic phase but dispersed when heated into the isotropic phase. Loudet *et al.* [10] dispersed silicone oil in a liquid-crystalline fluid

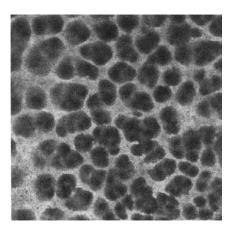


FIGURE 2 LSCM image of network structure with PMMA particles of radius 390 nm and volume fraction 5%.

and the mixture did not separate fully into two phases but self-organised into aligned chains of monodisperse colloidal droplets. In all cases, however, the density of particles is low, often vanishingly small.

In our case, we have succeeded in mixing a significant amount of particles (0.5 to 20 volume percent) with a nematic liquid crystal without macroscopic phase separation. The particles in fact form a three-dimensional network structure in the nematic liquid crystal (see Figure 2), conferring mechanical rigidity [4,6]. Upon crossing the isotropic to nematic phase transition of the LC walls of particles which enclose regions of liquid crystal were formed. With this structure the composite has a surprisingly high storage modulus of up to $10^7 \mathrm{Pa}$, showing the samples ability to be rigid and self-supporting over a wide temperature range. It turns out that the orientation of the liquid crystal molecules in this material can still be electrically switched [11].

In this paper we focus on the kinetics of the phase transition. Using Laser Scanning Confocal Microscopy (LSCM) we image the motion of single particles while the network structure is being formed. A microscope image of this structure is shown in Figure 2, where the liquid crystal appears as black and the PMMA as white dots. We gather detailed information on the kinetics of the transition, examining the particle motion, distribution of LC cell sizes and the wall formation during the process of the transition. We examine the reversibility of the network formation using microscopy and calorimetric measurements, and we discuss why the liquid crystal component of this novel composite appears to have more or less the properties it has in the pure state.

2. METHODOLOGY

Sample Preparation

We focus on mixtures of the model liquid crystal 5CB[†] (4-n-pentyl-4-cvanobiphenyl) and nearly monodisperse PMMA (polymethylmethacrylate) spheres[‡], with a particle radius of between 100 and 700 nm. A layer of poly-12-hydroxystearic acid (PHSA) is grafted onto the particles to sterically stabilize them. See Figure 1. PMMA particles of radius 200 nm, 390 nm and 500 nm are chemically bound with a fluorophore to enable observation by LSCM. The particles are cleaned by dispersing in hexane and then heated to evaporate off the solvent. The particles are dried for ~ 4 days at $\sim 65^{\circ}$ C in a vacuum oven. The particles are then dispersed in isotropic 5CB at approximately 45°C using a heated magnetic stirrer within an oven, and the samples were maintained in this way. At this temperature 5CB is in isotropic phase (isotropic-nematic transition temperature $T_{IN} \approx$ 35.3°C). Composites are made with particle volume fractions of $0.5\% \le$ $\phi \leq 15\%$. Prior to experimental use the sample is mixed on a whirl mixer, at $T \approx 45^{\circ}$ C, for approximately 30 minutes. Effort is taken to ensure that all samples are loaded into the respective apparatus in the isotropic state.

Microscopy Measurements

A Nikon TE300 inverted microscope is used with a Biorad Radiance 2000 laser scanning confocal microscopy unit [13]. A \times 100 oil-immersion objective with numerical aperture 1.4 is used. The confocal microscope is unique in that it allows a 'slice' within the bulk of the material to be imaged. An aperture in front of the objective filters out any out of focus rays from other planes or any scattered light, enabling a larger signal to noise ratio to be obtained. The excitation and emission wavelengths of the fluorophore are 488 nm and 520 nm respectively. With this optical setup and fluorophore the depth of field is 0.52 µm thick. For the 390 nm and 500 nm particles this means that only one layer of particles were imaged at any time. The typical time between successive images is 1.2s. Samples are prepared for observation within an oven (T $\approx 45^{\circ}$ C), by filling a 100 μ m capillary with the composite. Samples are prepared with particle volume fractions of $2\% \le \phi \le 10\%$. An enclosed heating stage enables the sample to be imaged under controlled temperatures and for the initial network formation to be observed under decreasing temperature. The sample is in contact with the objective oil where the sample temperature would ideally be measured and

[†]The 5CB was supplied by Merck.

[‡]The PMMA particles were made by Andy Schofield at The University of Edinburgh [12].

as a result temperature differences between the sample temperature and the measured temperature may exist.

DSC Measurements

A VP-DSC Microcalorimeter (Microcal Inc.) is used to make differential scanning calorimetry measurements [14]. The calorimeter used has two coin shaped cells fixed in place of $0.5\,\mathrm{ml}$ each. They are loaded through narrow tubes. Differences in the heat capacity between the reference cell and the sample cell are measured, where the reference cell is filled with water. The cells are cleaned with toluene and then ethanol and allowed to air dry before loading the sample. The sample is loaded into the preheated cell using a syringe that is also heated to approximately $55^{\circ}\mathrm{C}$. Generally measurements are made with a cooling rate of $6\,\mathrm{K/h}$ and a heating rate of $30\,\mathrm{K/h}$.

3. EXPERIMENTAL RESULTS

3.1. LSCM: Dynamics of the Phase Transition

A sample of PMMA particles of radius 390 nm and concentration 10% is prepared in the isotropic phase, homogenized for several days, loaded into the capillary and placed in the enclosed temperature stage. The temperature of the sample is decreased slowly in order to image the initial network formation. Images from this sequence are shown in Figures 3a–3i and show the progress of the network formation. The particles appear as white dots and the liquid crystal as black.

Figure 3a shows an image of the mixture in the isotropic phase in which individual particles can be resolved. No particle aggregation is observed and the particles exhibit Brownian motion with a diffusion constant of $D \approx 10^{-14} \,\mathrm{m}^2 \mathrm{s}^{-1}$. After passing the isotropic to nematic phase transition regions of liquid crystal begin to nucleate, expelling the PMMA particles. As can be seen in Figure 3b these regions are generally spherical. These regions of liquid crystal then begin to grow and locally the particle density increases as can be observed in Figure 3c. The particles form walls and the thin walls break during coalescence. Figure 3d and 3e show these growing regions coalescing and becoming elongated. From the original nucleation to this stage only requires approximately 10 seconds. Upon reaching this stage the kinetics of the phase separation are dominated by a slower process. The walls begin to compactify until motion becomes arrested and particles remain frozen in position. This is observed in images 3e and 3f and occurs over a period of approximately 20 seconds, however these timescales will depend upon the cooling rate. The thickness of the walls is ~ 10 particle diameters but they are much thicker close to nodes, and there

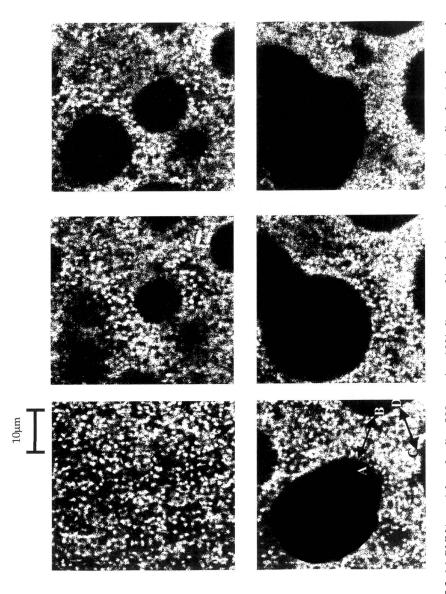


FIGURE 3 (a) PMMA particles of radius 390 nm, $\phi = 10\%$ dispersed in the isotropic phase. A scaling bar is shown at the top of the (d) Image taken at t = 3.6 seconds, showing coalescence of LC domains. AB and CD denote positions where the thickness of the image and remains constant throughout the image. White dots: colloids. (b) Image of sample just below the transition temperature showing the onset of phase separation, t = 0. (c) Image taken at t = 1.2 seconds, showing the growth of the liquid crystal domains. walls is measured. (e) Image taken at t = 9.6 seconds. (f) Image taken at t = 33.6 seconds.

is a broad distribution in the size of the liquid crystal domains. The smallest domains being the regions between the closest packed particles in the walls that are ~ 1 particle diameter. There are larger LC domains within the walls (see Fig. 3f) and there are the large domains enclosed by the walls that are as big as $\sim 50-100$ particle diameters.

To obtain detailed information on the compactification of the walls imaging software is used to measure the thickness of the walls at AB and CD (see Fig. 3d) in successive images. This is carried out from the stage where the original walls are formed until the motion of the network appears frozen. See Figure 4. There is a fluctuation in the thickness of the walls when two nearby domains coalesce. Figure 4 reveals that there is a fast decrease in the thickness within the first 10 seconds, a slower decrease over the next 10 seconds and it is almost constant thereafter.

The network formation appears reversible up to a point. Imaging has shown that upon heating the outermost particles become free until the network is broken up, however successive images taken at a higher magnification reveal that clusters of particles remain.

3.2. Differential Scanning Calorimetry

It has been suggested that impurities can cause the phase transition temperature to shift because of the disorder induced in the director. If the PMMA particles are classed as impurities it might be expected that they cause the transition temperature to decrease. As DSC is a sensitive method

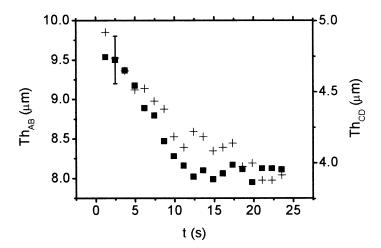


FIGURE 4 The thickness of the walls AB (\spadesuit) and CD (+) (see Figure 3f) as the network is formed. R = 390 nm, $\phi = 10\%$.

66 Cleaver et al.

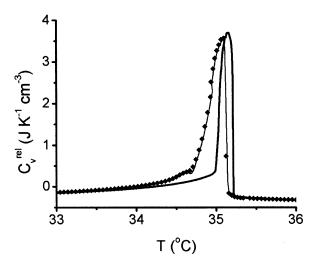


FIGURE 5 1st (\blacksquare) and 4th (\spadesuit) cooling with particles of radius R = 460 nm and concentration $\phi = 10\%$. Pure 5CB (\blacksquare) Cooling rate = 4.5 K/h from 45°C.

to detect the transition temperature cooling scans of 5CB and composite are run. The thick solid black line in Figure 5 represents the temperature dependence of the specific heat of pure 5CB. The specific heat curve of the pure 5CB remains constant until $\sim 35.2^{\circ}$ C, where there is a sharp peak of $\sim 0.2^{\circ}$ width. The thin solid black line represents the 1st cooling and the \rightarrow scatter graph shows the 4th cooling of the composite without changing sample material. The specific heat curves of the 1st and 4th cooling are constant until there is a sharp peak at $\sim 35.1^{\circ}$ C of width $\sim 0.3^{\circ}$. The similar position and shape of the 1st and 4th cooling curves demonstrates the reproducibility of the phase transition upon successive cooling. The position of the composite curves relative to the pure 5CB curve shows a transition temperature difference of only $\sim 0.1^{\circ}$ C. Thus concluding that the transition temperature of the liquid crystal is almost unaffected by the presence of the PMMA particles and strongly suggesting that the liquid crystal behaves as it does in the bulk.

4. CONCLUSIONS

We have carried out a series of measurements on the formation of colloidal networks in colloid liquid crystal composites using a combination

These DSC measurements do not support findings published by Petrov and Terentjev [15]. These differences will be discussed in a forthcoming publication

of laser-scanning confocal imaging and DSC techniques. We have obtained detailed (single particle resolution) data showing the formation of the network upon cooling from the initial isotropic dispersion. The sequence of images show the mechanism for network formation and suggest that it is a continuous process in which nematic domains grow, expelling particles to form a complex three dimensional network. Calorimetry measureisotropic-nematic transition show ments through the that thermodynamic aspects of the isotropic-nematic transition appear similar to those of the bulk liquid crystal material. This is likely, since imaging reveals the smallest liquid crystal domains in these composites to be typically ~ 100 nm which is an order of magnitude larger than the correlation length of the liquid crystal molecule ~10nm. This finding is consistent with the observation that the LC switching times in the composites are similar to those of the pure liquid crystal.

REFERENCES

- de Gennes P. G. & Prost J. (1993). The Physics of Liquid Crystals, Oxford University Press, New York.
- [2] Demus, D. Goodby, J. Gray, G. W. Spiess, H. W., & Vill, V. (1998). Handbook of Liquid Crystals, Vol 1: Fundamentals, Wiley-VCH, Weinheim.
- [3] Gregory Philip Crawford, & Slobodan Zumer. (1996). Liquid Crystals in Complex Geometries, Taylor and Francis, London.
- [4] Meeker, S. P. Poon, W. C. K., Crain, J., & Terentjev, E. M. (2000). Colloid-liquid-crystal composites: An unusual soft solid, *Phys. Rev E Rapid Communications*, 61, R6083– R6086.
- [5] Anderson, V. J., Terentjev, E. M., Meeker, S. P., Crain, J., & Poon, W. C. K. (2001). Cellular solid behavior of liquid crystal colloids 1. Phase separation and morphology, Eur. Phys., J. E, 4, 11–20.
- [6] Anderson, V. J. & Terentjev, E. M. (2001). Cellular solid behavior of liquid crystal colloids 2. Mechanical properties. Eur. Phys. J. E, 4, 21–28.
- [7] Yuedong Gu, & Nicholas L. Abbot. (2000). Observation of Saturn-Ring Defects around Solid Microspheres in Nematic Liquid Crystals. *Physical Review Letters*, 22, 4719–4722.
- [8] Terentjev, E. M. (1995). Disclination loops, standing alone and around solid particles, in nematic liquid crystals. *Physical Review E*, 51, 2, 1330–1337.
- [9] Philippe Poulin, Holger Stark, Lubensky, T. C., & Weitz, D. A. (1997). Novel interactions in anisotropic fluids. Science, 275, 1770–1773.
- [10] Jean-Christophe Loudet, Phillipe Barois, and Philipe Poulin. (2000). Colloidal ordering from phase separation in a liquid-crystalline continuos phase. *Nature*, 407, 611–613.
- [11] Nakamura, H., Meeker, S. P., Taira, Y., Poon, W. C. K., & Crain, J. (2001). Optical properties of colloid liquid crystal composites. Mol. Cryst. and Lig. Cryst., 368, 167–174.
- [12] Antl, Goodwin, J. W., Hill, R. D., Ottewill, R. H., Owens, S. M., Papworth, S., & Waters, J. A. (1986). The preparation of poly(methylmethacrylate) latices in non-aqueous media. *Colloid and Surfaces*, 17, 67–78.

68 Cleaver et al.

- [13] Sheppard, C. J. R. & Shotton, D. M. (1997). Confocal Laser Scanning Microscopy, BIOS Scientific Publishers, Oxford.
- [14] Plotnikov, V. V., Brandts, J. M., Lin, L. N., & Brandts, J. F. (1997). A new ultrasensitive calorimeter. *Analytical Biochemistry*, 250, 237–244
- [15] Peter G. Petrov & Eugene M. Terentjev. (2001). Formation of Cellular Solid in Liquid Crystal Colloids. *Langmuir*, 17, 2942–2949.